
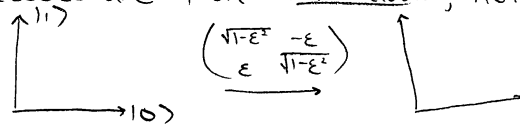


Problems for implementing a quantum computer:

- Quantum entangled states are fragile
 eg., measuring a single qubit of $|0^n\rangle + |1^n\rangle$, the cat state, collapses it entirely

Similarly  interacting with a single photon and then tracing it out

- Error processes are often continuous, not discrete



and small errors can accumulate

- There is no generic method for duplicating quantum information

No cloning theorem: There is no linear map (let alone an isometry) that maps $|ψ\rangle \mapsto |ψ\rangle \otimes |ψ\rangle$ for all $|ψ\rangle$

Proof: If $|0\rangle \mapsto |0\rangle \otimes |0\rangle$ and $|1\rangle \mapsto |1\rangle \otimes |1\rangle$, then

linearity requires $\alpha|0\rangle + \beta|1\rangle \mapsto \alpha|00\rangle + \beta|11\rangle$

$$\neq (\alpha|0\rangle + \beta|1\rangle) \otimes (\alpha|0\rangle + \beta|1\rangle) \quad \square$$

(can't make backups!)

- We want to compute on the data, not just protect it in memory
 Computations themselves need to be protected / "fault tolerant"

- Noise rates in controllable quantum systems are usually quite high

eg. [1009.2267] compares some systems (environmentally accessible)

superconducting charge qubit 1% noise for one-qubit ops

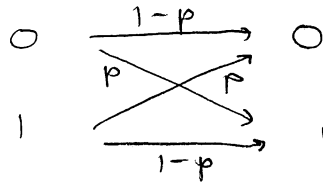
classical devices have lower noise: 10% noise for two-qubit ops

transistor noise rate 10^{-16} ?

neuron noise rate ?

Codes for data memory or transmission

Binary symmetric channel / bit-flip noise



protect against using 3-bit repetition code (simplest example)

encode transmit decode
 $0 \rightarrow 000$ \longrightarrow $abc \rightarrow \text{majority}(a,b,c)$
 $1 \rightarrow 111$

$$\begin{aligned} P[\text{erroneous transmission}] &= 3p^2(1-p) + p^3 \\ &= 3p^2 - 2p^3 \\ &< p \text{ for } p < 1/2 \quad \checkmark \end{aligned}$$

Quantum bit-flip channel (p)

With prob. $1-p$ apply I , and with prob. p apply $X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$.
 $E(p) = (1-p)\rho + pX\rho X$

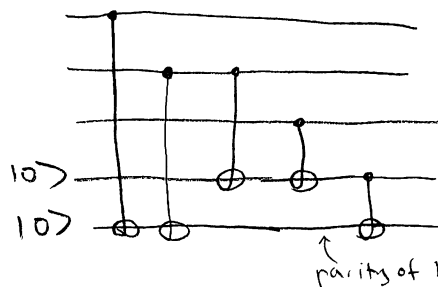
The same code can be used, with care.

encoding

$$\alpha|0\rangle + \beta|1\rangle \longrightarrow \alpha|1000\rangle + \beta|1111\rangle$$

eg. if there is an error on the third qubit
 transmission $\longrightarrow \alpha|1001\rangle + \beta|1110\rangle$

Decoding: - cannot measure and apply majority
 $\longrightarrow \begin{pmatrix} |\alpha|^2 & 0 \\ 0 & |\beta|^2 \end{pmatrix}$ collapsed superposition

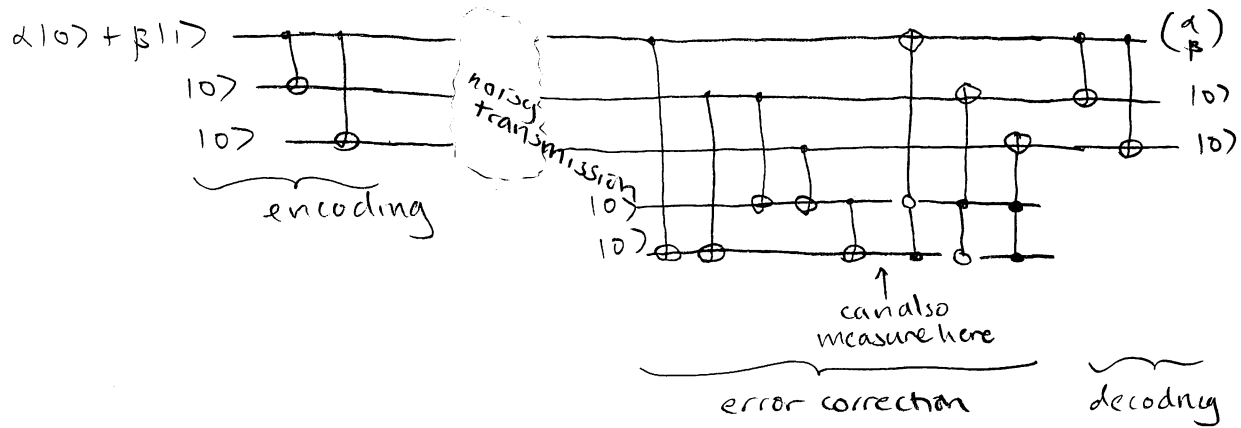


$$\begin{aligned} (\alpha|1001\rangle + \beta|1110\rangle)|00\rangle &\longrightarrow \alpha|1001\rangle|11\rangle + \beta|1110\rangle|11\rangle \\ &= (\alpha|1001\rangle + \beta|1110\rangle)|11\rangle \end{aligned}$$

"syndrome"

Classical state	Syndrome
000	00
001	11
010	10
011	01
100	01
101	10
110	11
111	00

- no errors \Rightarrow syndrome is 00
 - X on qubit $j \Rightarrow$ syndrome is j in binary
- \Rightarrow apply a correction conditioned on the syndrome

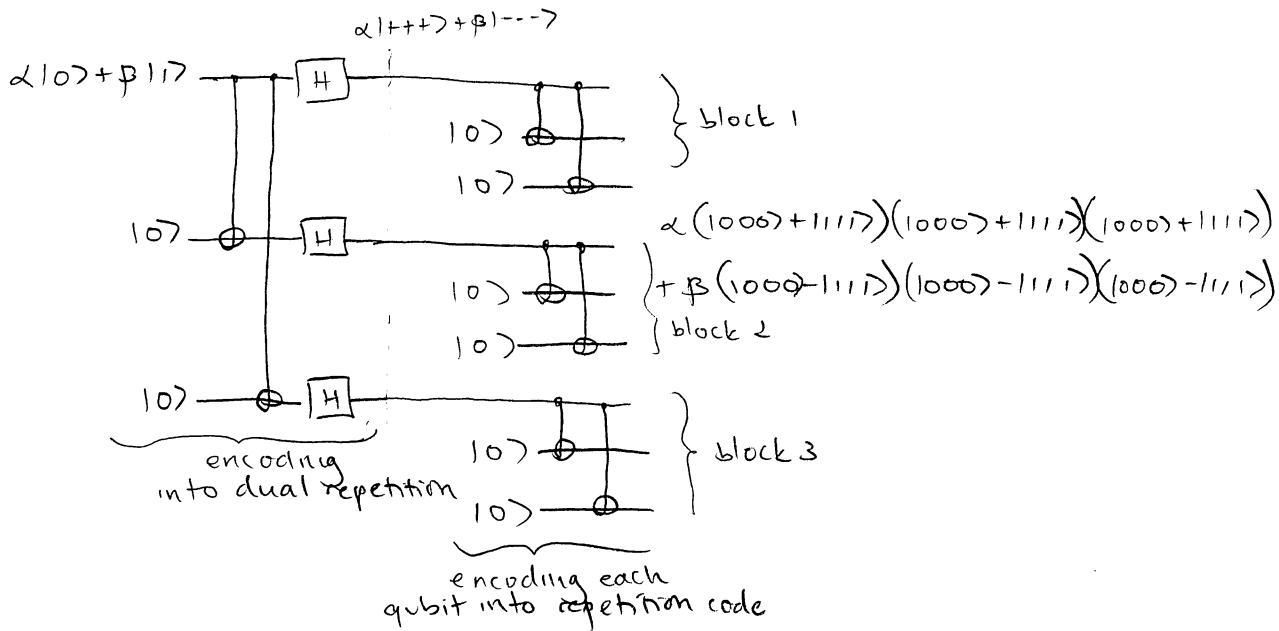


Correcting both phase and bit-flip errors

- use both the previous codes, i.e., "concatenate" them

since $Z_3(\alpha|1000\rangle + \beta|1111\rangle) = \alpha|1000\rangle - \beta|1111\rangle$

a phase-flip on the encoded states $|1000\rangle, |1111\rangle$, we can protect against it with another layer of coding



Shor's 9-qubit code:

Can correct up to one X error in each block, and up to one Z error total.

⇒ can protect against any of the errors X, Z or Y = iXZ on a single qubit

(Note: Can correct errors without decoding)

Example: Apply XZ to qubit 4 of the state...

Phase-flip channel ("dephasing", " T_2 ")

$$\mathcal{E}(\rho) = (1-p)\rho + pZ\rho Z \quad Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$Z|b\rangle = (-1)^b|b\rangle$$

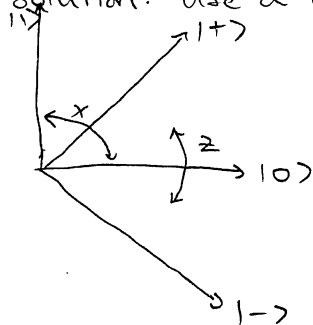
Observe: The above repetition code makes phase errors worse!

a Z on any one qubit will map

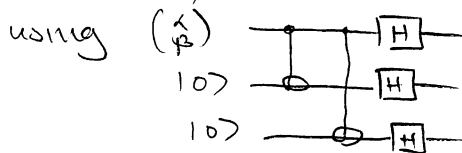
$$\alpha|100\rangle + \beta|111\rangle \mapsto \alpha|100\rangle - \beta|111\rangle$$

i.e., apply a logical phase flip.

Simple solution: Use a repetition code in the dual basis



$$\text{encode } \alpha|10\rangle + \beta|11\rangle \mapsto \alpha|++\rangle + \beta|--\rangle$$



phase flips now act the same as bit flips did before

\Rightarrow Bob can apply Hadamards and then the same correction and decoding circuits

General quantum errors

Key point: It suffices for a quantum code to correct against the Pauli errors, because I, X, Y, Z form a basis for all 2×2 matrices.

Example: Let U be a 2×2 unitary

$$U = aI + bX + cY + dZ$$

Let $|\mathcal{F}\rangle$ be an encoded quantum state

$$U^{(j)}|\mathcal{F}\rangle = a|\mathcal{F}\rangle + bX^{(j)}|\mathcal{F}\rangle + cY^{(j)}|\mathcal{F}\rangle + dZ^{(j)}|\mathcal{F}\rangle$$

$$\begin{array}{l} \xrightarrow{\text{syndrome}} \\ \text{extraction} \end{array} \quad \begin{array}{l} a|\mathcal{F}\rangle \otimes |I \text{ syndrome}\rangle \\ + bX^{(j)}|\mathcal{F}\rangle \otimes |X^{(j)} \text{ syndrome}\rangle \\ + cY^{(j)}|\mathcal{F}\rangle \otimes |Y^{(j)} \text{ syndrome}\rangle \\ + dZ^{(j)}|\mathcal{F}\rangle \otimes |Z^{(j)} \text{ syndrome}\rangle \end{array}$$

$$\xrightarrow{\text{correction}} \quad |\mathcal{F}\rangle \otimes \left(\begin{array}{l} a|I \text{ syndrome}\rangle + b|X^{(j)} \text{ syndrome}\rangle \\ + c|Y^{(j)} \text{ syndrome}\rangle + d|Z^{(j)} \text{ syndrome}\rangle \end{array} \right)$$

Having recovered the codeword $|\mathcal{F}\rangle$, we can discard the extra ancillas.

More generally, consider an arbitrary quantum map (CPTP map)

$$\mathcal{E}^{(j)}(\rho) = \sum_{k=1}^m A_k^{(j)} \rho (A_k^{(j)})^\dagger$$

Each Kraus operator can be expanded in the Pauli basis

$$A_k = a_k I + b_k X + c_k Y + d_k Z$$

For notational simplicity let $\sigma_0 = I, \sigma_1 = X, \sigma_2 = Y, \sigma_3 = Z,$

$$\Rightarrow \mathcal{E}^{(j)}(|\mathcal{F}\rangle\langle\mathcal{F}|) = \sum_{k=1}^m \sum_{p=0}^3 \sum_{r=0}^3 a_{k,r} \sigma_r^{(j)} |\mathcal{F}\rangle\langle\mathcal{F}| \sigma_p^{(j)} \overline{a_{k,r}}$$

$$\xrightarrow{\text{syndrome}} \sum_{k=1}^m \sum_{p=0}^3 \sum_{r=0}^3 a_{k,r} \overline{a_{k,r}} \left(\sigma_p^{(j)} |\mathcal{F}\rangle\langle\mathcal{F}| \sigma_r^{(j)} \right) \otimes |\sigma_p^{(j)} \text{ syn.}\rangle \langle \sigma_r^{(j)} \text{ syn.}|$$

$$\xrightarrow{\text{correction}} \quad |\mathcal{F}\rangle\langle\mathcal{F}| \otimes (\text{stuff})$$

Observe: Measuring the syndrome would give

$$\sum_k \sum_r |a_{k,r}|^2 \sigma_p^{(j)} |\mathcal{F}\rangle\langle\mathcal{F}| \sigma_r^{(j)} \otimes |\sigma_p^{(j)} \text{ syn.}\rangle \langle \sigma_r^{(j)} \text{ syn.}|$$

→ thus it projects or collapses the arbitrary (perhaps continuous) error into just one of four discrete errors.

Codes that correct discrete I, X, Y, Z errors also correct continuous errors, eg $\begin{pmatrix} \cos\theta & 0 \\ \sin\theta & \cos\theta \end{pmatrix}$, automatically, for free!