

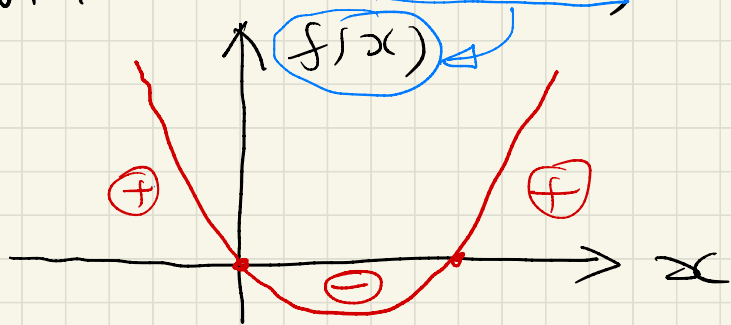

Lecture 15

03/08/23

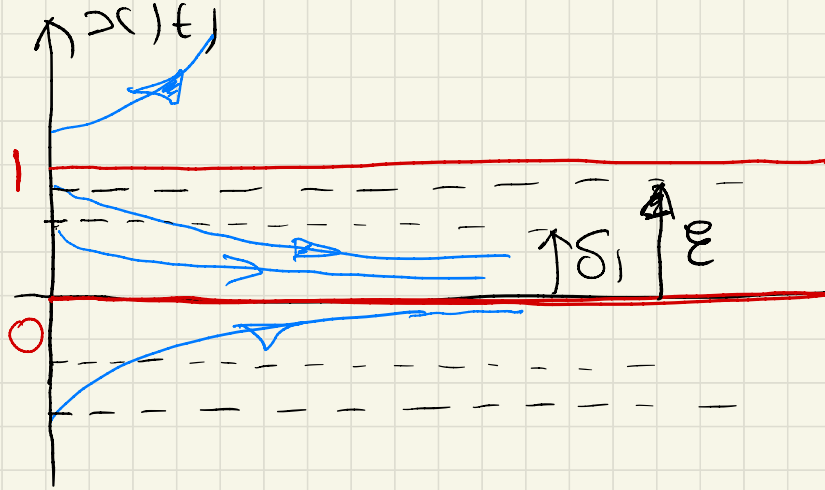
Today:

- Use of stability definitions (examples)
- Lyapunov-based method for stability analysis

Ex: $\dot{x} = x(x-1) \Rightarrow$ 2 eq. points
 $\bar{x} = 0; \bar{x} = 1$



①



We can use stability definition to directly check stability

properties of eq.

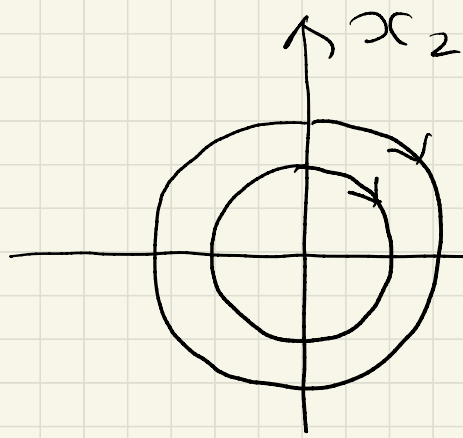
points in situations where we can determine solution to $\dot{x} = f(x)$.

In this example, $\bar{x} = 1$ is unstable
 $\bar{x} = 0$: LAS (locally asymptotically stable)

Ex: Harmonic oscillator

$$\begin{bmatrix} \dot{x}_1 \\ \dot{x}_2 \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix}; \quad \ddot{y} + y = 0$$

$$x_1 = y; \quad x_2 = \dot{y} \quad (2)$$



$\bar{x}=0$; stable but
not asymptotically
stable.

~~$(x(t) \xrightarrow{t \rightarrow \infty} 0)$~~

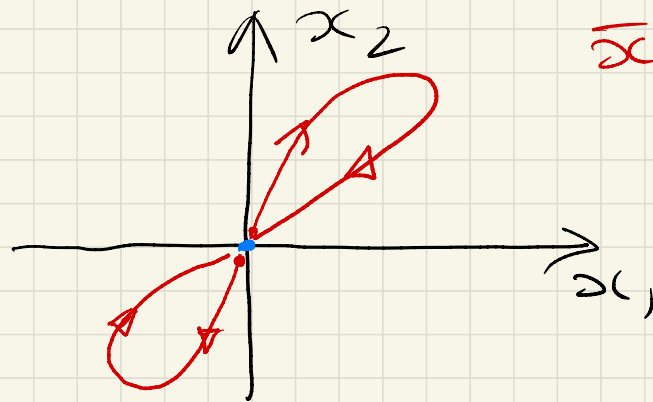
For linear systems,

$\lim_{t \rightarrow \infty} \|x(t)\| = 0 \Rightarrow \bar{x}=0$ is stable
eq. point.

In nonlinear case, attractiveness
of $\bar{x}=0$ (i.e., $\lim_{t \rightarrow \infty} \|x(t)\| = 0$)

DOES NOT imply stability of $\bar{x}=0$. (3)

Ex: Homoclinic orbit



$\bar{x}=0$ is not stable but

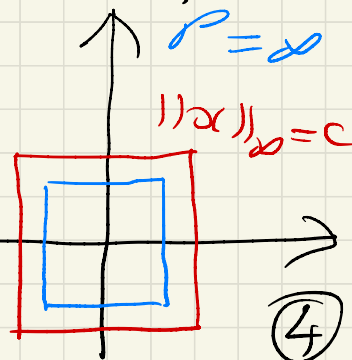
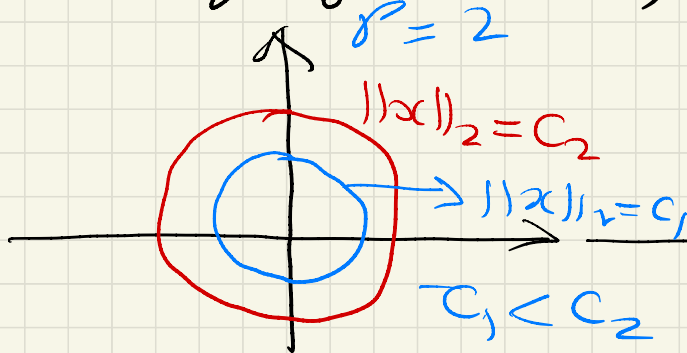
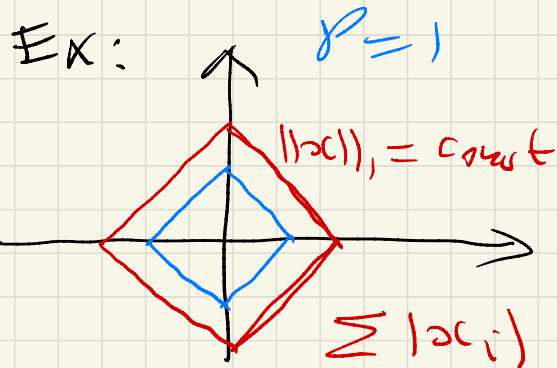
$$\lim_{t \rightarrow \infty} \|x(t)\| = 0$$

$t \rightarrow \infty$

See Khalid for details

Note on norms: $\|\cdot\|$ in stability

definition is any p -norm; $p \geq 1$



Lyapunov-based method for stability of eq. points of $\dot{x} = f(x)$

Motivating example (pendulum)

$$\begin{bmatrix} \dot{x}_1 \\ \dot{x}_2 \end{bmatrix} = \begin{bmatrix} x_2 \\ -a \sin x_1 - b x_2 \end{bmatrix}$$

$x_1 = \theta$
 $x_2 = \dot{\theta}$

$$\ddot{\theta} + b \cdot \dot{\theta} + a \cdot \sin \theta = 0$$

Eq. points: $\bar{x} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$ "down" or $\begin{bmatrix} \pi \\ 0 \end{bmatrix}$ "up"

Neither eq. point CAN be GAS

(because of the existence of another eq. point !!!)

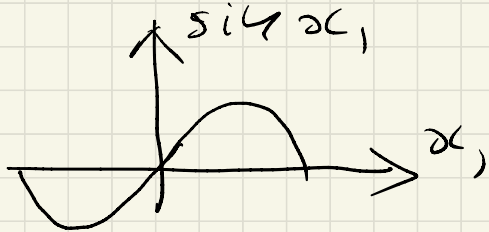
(5)

Energy:

$$E(x_1, x_2) = \underbrace{c_1 \cdot \int_0^{x_1} \sin \xi \, d\xi}_{\text{potential}} + \underbrace{c_2 \cdot \frac{1}{2} x_2^2}_{\text{kinetic}}$$

$$\rightarrow E(0, 0) = 0 \quad c_1(1 - \cos x_1)$$

$E(x_1, x_2) > 0$ on domain D around the origin



$$x_1, \sin x_1 > 0, \quad x_1 \in (-\pi, \pi)$$

Objective: study $\frac{d}{dt} E$ along solutions to $\ddot{x} = f(x)$.

$$E(t) = E(x_1(t), x_2(t))$$

$$\frac{dE}{dt} = \left[\nabla E \right]^T \cdot \frac{dx}{dt} = \underbrace{\left[\frac{\partial E}{\partial x_1} \quad \frac{\partial E}{\partial x_2} \right]}_{\left[\nabla E \right]^T} \cdot \begin{bmatrix} \dot{x}_1 \\ \dot{x}_2 \end{bmatrix}$$

$$= \left[\frac{\partial E}{\partial x_1} \quad \frac{\partial E}{\partial x_2} \right] \cdot \begin{bmatrix} f_1 \\ f_2 \end{bmatrix} = f_1 \frac{\partial E}{\partial x_1} + f_2 \frac{\partial E}{\partial x_2}$$

$$= \sum_{i=1}^2 f_i(x_1, x_2) \cdot \frac{\partial E(x_1, x_2)}{\partial x_i}$$

Holds for any 2nd order system

For pendulum:

$$\frac{\partial E}{\partial x_1} = C_1 \cdot \sin x_1 ; \quad \frac{\partial E}{\partial x_2} = C_2 \cdot x_2$$

$$f_1 = x_2 ; \quad f_2 = -m g \sin x_1 - b x_2 \quad \Rightarrow \quad \textcircled{7}$$

$$\begin{aligned}\frac{dE}{dt} &= C_1 \cdot \underline{x_2 \cdot \sinh x_1} - C_2 \underline{x_2} \left[\underline{a \sinh x_1} + \underline{b x_2} \right] \\ &= \underbrace{[C_1 - C_2 \cdot a]}_{\text{sign-indefinite}} \cdot \underbrace{x_2 \sinh x_1}_{\text{sign-indefinite}} - C_2 \cdot b \cdot \underbrace{x_2^2}_{\text{always negative}}\end{aligned}$$



sign-indefinite
(i.e. can be both positive
and negative around
 $\bar{x}=0$)

(always
negative)

The best we can do: choose C_1 and C_2
to cancel it.

E.g. $C_2 = 1$; $C_1 = a$ works !!!

With this choice $\Rightarrow \frac{dE}{dt} = -b \cdot x_2^2 \leq 0$ (8)

Note: $\frac{dE}{dt} = 0 \cdot x_1^2 - 2 \cdot x_2^2 \leq 0$ or ≥ 0

$x_1 \neq 0$ and $x_2 = 0 \Rightarrow \frac{dE}{dt} = 0$

$\Rightarrow E(t)$ is a non-increasing function of time.

Allows us to conclude stability (in the sense of Lyapunov) of $\bar{x} = 0$.

(but not local asymptotic stability): further analysis or different choice of energy-like function required. ③

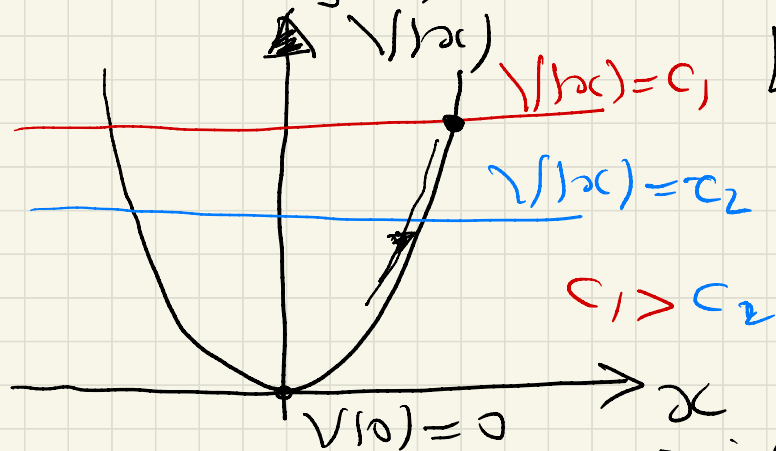
If $\gamma = 0 \Rightarrow$ no viscous damping

$$\frac{dE}{dt} = 0 \Rightarrow E(t) = E(t_0) = \text{const.}$$

Conservative system [Energy does not increase or decrease along the solutions]
(Same as harmonic oscillator).
Linearization of pendulum
with $\gamma = 0$ around $\bar{x} = 0$.

Lagrange-based method for statistical analysis of equilibria generalizes the above argument to a broader class of energy-like functions for $\bar{x} = f(x)$.

Ex: $\dot{x} = f(x)$ with $x(t) \in \mathbb{R}$



Lyapunov function
candidate

$V(x)$ [instead of $E(x)$]

- represents an initial condition

Want to figure out if V increases;
decreases or
stays constant
along the solutions of
 $\dot{x} = f(x)$


$$E_x: \quad \partial C = x(x-1)$$

$$V(x) = \frac{1}{2} x^2 \quad \left(V(0) = 0; \quad V(x) > 0 \text{ for all } x \neq 0 \right)$$

$$V' = \frac{\partial V}{\partial x} \cdot \partial C = x \cdot x(x-1) =$$

$$= - \underbrace{(1-x)}_{f(x)} x^2$$

$f(x)$

want it 

to be > 0 for $x \neq 0$

$$\text{For } x < 1 \Rightarrow \frac{\partial V}{\partial t} < 0 \text{ for all } x \neq 0$$

$$\frac{dV}{dt} = 0 \text{ for } x = 0 \quad \left[\text{by construction} \right]$$

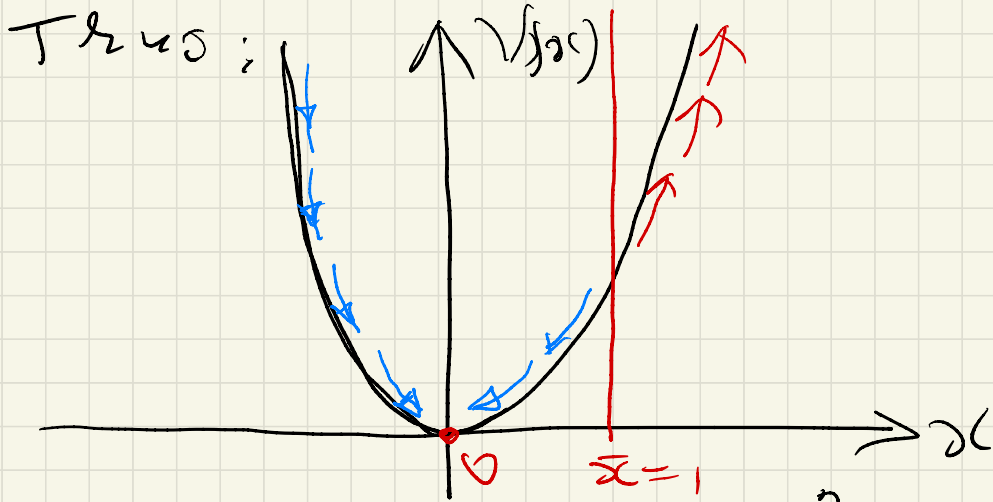
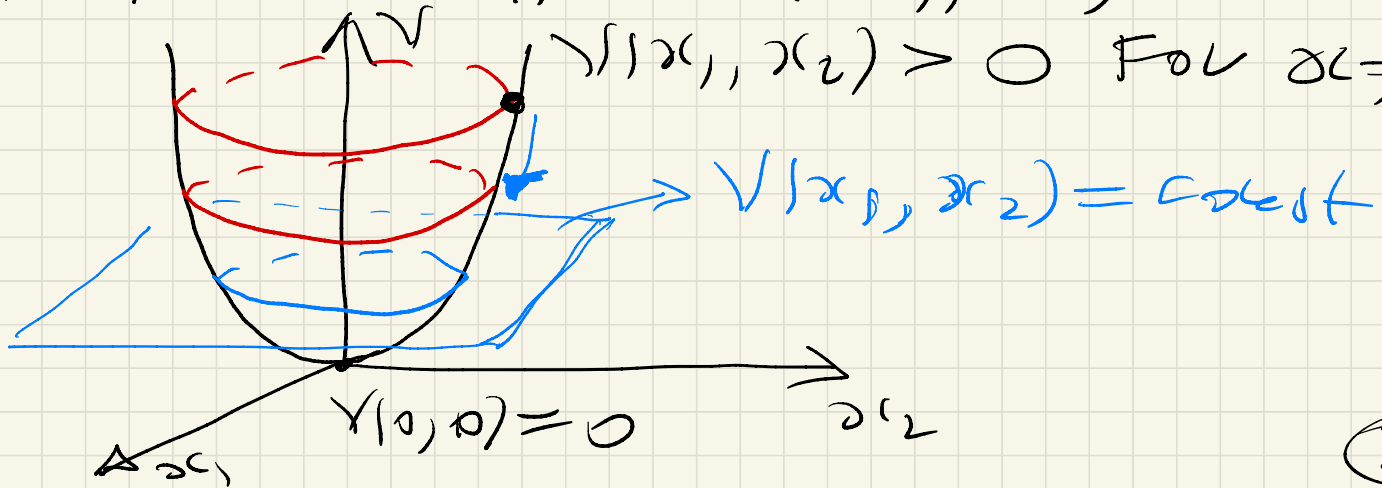


Illustration in \mathbb{R}^2 $V(x_1, x_2)$



Thm: Let D be an open connected subset of \mathbb{R}^n that contains the eq. point $\bar{x} = 0$ of $\dot{x} = f(x)$ [$f(0) = 0$].
 1° If there is a continuously differentiable function

$$V: D \rightarrow \mathbb{R}$$

such that V : locally positive definite

$$a) V(0) = 0$$

$$b) V(x) > 0 \text{ for all } x \in D \setminus \{0\}$$

and $\frac{dV}{dt} = [\nabla V]^T \cdot f(x) \leq 0$ for all $x \in D$ [\dot{V} : locally negative ^{semi-definite}]

then $\bar{u} = 0$ is stable in the sense of Lyapunov.